

# Kinetics of fluid demixing in complex plasmas

## Morphological Data Analysis using Minkowski Tensors

Alexander Böbel and Christoph R  th

DLR German Aerospace Center  
Research Group Complex Plasmas

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Knowledge for Tomorrow



# Outline

1. Minkowski Valuations
2. Elementary Applications
3. Simulated demixing of a binary complex plasma
4. Outlook



# Brezen!



Fig. 1 : Above: a Bavarian brezen ('pretzel')

$$W_2(K) = \#(\text{connected regions}) - \#(\text{holes}) \quad (1)$$



# Minkowski Functionals

For a body  $K$  with a smooth boundary contour  $\partial K$  embedded in  $D$ -dim euclidean space the  $D + 1$  Minkowski functionals are defined as:

$$\rightarrow W_0(K) = \int_K d^D r$$

$$\rightarrow W_\nu(K) = \int_{\partial K} G_\nu(r) d^{D-1} r \quad , \quad 1 \leq \nu \leq D$$

$G_\nu(r)$  are the elementary symmetric polynomials of the local principal curvatures.

in 2d:

$$\rightarrow W_0(K) = \int_K d^2 r \quad \propto \#(\text{pixels}) \quad (\text{area})$$

$$\rightarrow W_1(K) = \int_{\partial K} dr \quad \propto 2\#(\text{edges}) - 4\#(\text{squares}) \quad (\text{circumference})$$

$$\rightarrow W_2(K) = \int_{\partial K} \kappa(r) dr \quad \propto 2\#(\text{edges}) - 4\#(\text{squares}) \quad (\text{euler characteristic})$$

similar in 3d:

$$\rightarrow W_0(K) = \int_K d^3 r \quad (\text{volume})$$

$$\rightarrow W_1(K) = \int_{\partial K} d^2 r \quad (\text{area})$$

$$\rightarrow W_2(K) = \int_{\partial K} \kappa_1 + \kappa_2 d^2 r \quad (\text{integrated mean curvature})$$

$$\rightarrow W_3(K) = \int_{\partial K} \kappa_1 \kappa_2 d^2 r \quad (\text{euler characteristic})$$





# Properties and applications

MFs are motion invariant, additive, continuous. They form a complete family of morphological measures. They are sensitive to higher order correlations.

Applications:

- curvature energy of membranes
- order parameter in Turing patterns
- density functional theory for fluids (as hard balls or ellipsoids)
- test point distributions (find clusters, filaments, underlying pointprocess)
- non-gaussianity of CMB



# Abstraction to tensor valued valuations

In order to account also for directional properties it is natural to abstract the scalar valued MF to tensor valued quantities called MT:

## Definition

$$\rhd W_0^{a,0}(K) := \int_K d^D r \, \mathbf{r}^{\odot a} \quad , \quad (\nu = b = 0)$$

$$\rhd W_\nu^{a,b}(K) := 1/D \int_{\partial K} d^{D-1} r \, G_\nu(r) \, \mathbf{r}^{\odot a} \odot \mathbf{n}^{\odot b} \quad , \quad (\nu, b \neq 0)$$

Properties:

- $\rhd$  They are isometry covariant, i.e their behaviour under translation and rotation is given by:

$$W_\nu^{a,b}(K + \mathbf{t}) = \sum_{i=0}^a \binom{a}{i} \mathbf{t}^i W_\nu^{a-i,b}(K) \quad (2a)$$

$$W_\nu^{a,b}(\hat{O} K) = \hat{O}_{a+b} W_\nu^{a,b}(K) \quad (2b)$$

- $\rhd$  They are additive:  $W_\nu^{a,b}(K_1 \cup K_2) = W_\nu^{a,b}(K_1) + W_\nu^{a,b}(K_2) - W_\nu^{a,b}(K_1 \cap K_2)$

- $\rhd$  They are homogeneous:  $W_\nu^{a,b}(\lambda K) = \lambda^{3+a-\nu} W_\nu^{a,b}(K)$



# Completeness

Let  $K^d$  denote the family of all compact convex subsets of the Euclidean space  $R^d$ . Let  $L$  be a linear space.

## Definition

A function  $\phi : K^d \rightarrow L$  is called a valuation if

$$\phi(K_1 \cup K_2) + \phi(K_1 \cap K_2) = \phi(K_1) + \phi(K_2) \quad (3)$$

for any  $K_1, K_2 \in K^d$  such that  $K_1 \cup K_2 \in K^d$

## Theorem (Alesker 1999)

Let  $\phi : K^d \rightarrow L$  be a continuous translation- and  $SO(d)$ -covariant valuation. Then  $\phi$  has the form

$$\phi(K) = \sum_j c_j W_j(K) \quad (4)$$

where  $W_j(K)$  is the  $j^{\text{th}}$  Minkowski valuation, and  $c_j \in L$  are fixed uniquely defined vectors.

➤ every morphological measure is a linear combination of Minkowski valuations



# Isotropy measure $\beta$

For a body  $K$  and each Minkowski tensor  $W_{\nu}^{a,b}(K)$  an isotropy index can be defined as the ratio between smallest and largest eigenvalue of the  $D \times D$ -Matrix representing each Minkowski tensor.

## Definition

$$\beta_{\nu}^{a,b}(K) := \frac{\lambda_{\min} \left( W_{\nu}^{a,b}(K) \right)}{\lambda_{\max} \left( W_{\nu}^{a,b}(K) \right)} \quad (5)$$

The dimensionless isotropy index is a pure shape measure. It is invariant under isotropic scaling of  $K$ .

examples:

- sphere, circle       $\beta = 1$
- cube, square       $\beta = 1$
- box       $\beta = \text{shorter/longer edge}$

⇒ isotropy measure in the sense of elongation



Consider the simplest rank 4 MT:  $W_1^{04}(K) = 1/3 \cdot \int_{\partial K} d^2r \, \mathbf{n}(\mathbf{r}) \otimes \mathbf{n}(\mathbf{r}) \otimes \mathbf{n}(\mathbf{r}) \otimes \mathbf{n}(\mathbf{r})$

→ translation invariant

→ symmetric  $(W_1^{04})_{ijkl} = (W_1^{04})_{(ijkl)} \Rightarrow 15$  independent elements in 3d

Rewrite  $W_1^{04}$  in the Mehrabadi supermatrix notation as a  $6 \times 6$  matrix:

$$M = \begin{pmatrix} S_{xxxx} & S_{xxyy} & S_{xxzz} & \sqrt{2} S_{xxyz} & \sqrt{2} S_{xxxx} & \sqrt{2} S_{xxxxy} \\ S_{xxyy} & S_{yyyy} & S_{yyzz} & \sqrt{2} S_{yyyz} & \sqrt{2} S_{yyxz} & \sqrt{2} S_{yyxy} \\ S_{xxzz} & S_{yyzz} & S_{zzzz} & \sqrt{2} S_{zzyz} & \sqrt{2} S_{zzxz} & \sqrt{2} S_{zzxy} \\ \sqrt{2} S_{xxyz} & \sqrt{2} S_{yyyz} & \sqrt{2} S_{zzyz} & 2 S_{yzyz} & 2 S_{yzxz} & 2 S_{yzxy} \\ \sqrt{2} S_{xxxx} & \sqrt{2} S_{yyxz} & \sqrt{2} S_{zzxz} & 2 S_{yzxz} & 2 S_{xzxz} & 2 S_{xzxxy} \\ \sqrt{2} S_{xxxxy} & \sqrt{2} S_{yyxy} & \sqrt{2} S_{zzxy} & 2 S_{yzxy} & 2 S_{xyxz} & 2 S_{xyxy} \end{pmatrix}$$

with  $S = W_1^{04}(K)/W_1(K)$

It is possible to define a distance measure on the space of bodies:

## Definition

$$\Delta(K_1, K_2) := \left( \sum_{i=1}^6 (\zeta_i(K_1) - \zeta_i(K_2))^2 \right)^{1/2} \quad (7)$$

It is a pseudometric. It is symmetric, the triangle inequality holds, but:

$$\Delta(K_1, K_2) = 0 \Leftarrow K_1 = K_2$$



# Symmetry pseudometric $\Delta$

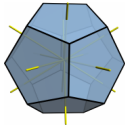
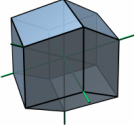
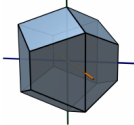
structure	ico	fcc	hcp	bcc	sc
type	(5,1)	(3,2,1)	(2,2,1,1)	(3,2,1)	(3,2,1)
$\zeta_1$	$\frac{1}{3}$	$\frac{1}{3}$	$\frac{1}{3}$	$\frac{1}{3}$	$\frac{1}{3}$
$\zeta_2$	$\frac{2}{15}$	$\frac{1}{6}$	$\frac{1}{6}$	$\frac{8-4/\sqrt{3}}{33}$	$\frac{1}{3}$
$\zeta_3$	$\frac{2}{15}$	$\frac{1}{6}$	$\frac{5}{36}$	$\frac{8-4/\sqrt{3}}{33}$	$\frac{1}{3}$
$\zeta_4$	$\frac{2}{15}$	$\frac{1}{6}$	$\frac{5}{36}$	$\frac{8-4/\sqrt{3}}{33}$	0
$\zeta_5$	$\frac{2}{15}$	$\frac{1}{12}$	$\frac{1}{9}$	$\frac{-1+2\sqrt{3}}{33}$	0
$\zeta_6$	$\frac{2}{15}$	$\frac{1}{12}$	$\frac{1}{9}$	$\frac{-1+2\sqrt{3}}{33}$	0
					

Fig. 2 : Above: eigenvalue tuple for ideal polyhedra, Kapfer 2011



# Ideal crystal with noise

Rank 2  $\beta$  MT analysis for ideal hcp, fcc, bcc data with gaussian noise. → **Isotropy index drops exponentially with increasing noise, independent of crystal type.**

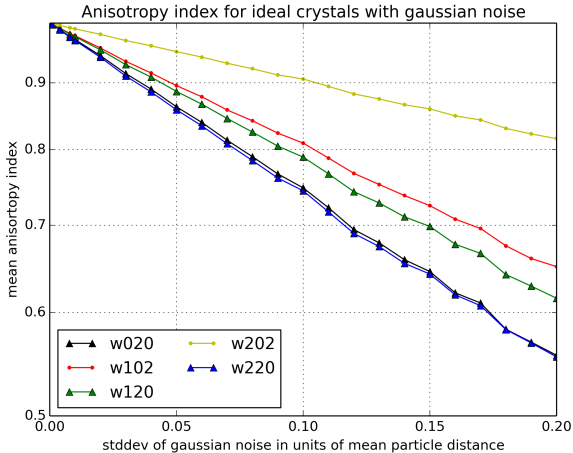
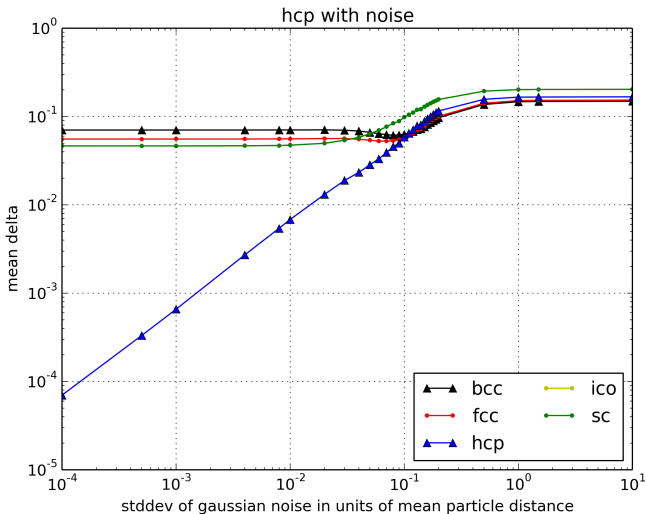


Fig. 3 : Above: Isotropy index drops exponentially with increasing noise.



# Ideal crystal with noise

Rank 4  $\Delta$  MT analysis for ideal hcp data with gaussian noise.





# Simulated demixing of a binary complex plasma

Plasma conditions:

- stationary isotropic highly collisional plasma
- Plasma production due to electron impact ionization
- Plasma losses due to three-body bulk recombination and ambipolar diffusion to the plasma boundaries.

The continuity and momentum equations for e.g. ions are:

$$\begin{aligned}\nabla(n_i v_i) &= \nu_i n_e - \nu_L n_i - \beta n_e n_i \\ (v_i \nabla) v_i &= -(e/m_i) \nabla \Phi - (u_{T_i}^2 / n_i) \nabla n_i - \nu_i v_i\end{aligned}\quad (8)$$

(here  $v_i$  and  $m_i$  are the velocity and mass of the ions,  $\nu_i$  the ionization frequency,  $\nu_L$  the characteristic frequency of ambipolar losses,  $\beta$  the recombination coefficient,  $\nu_i$  the characteristic frequency of ion-neutral collisions,  $\Phi$  the electrical potential, and  $u_{T_i}$  is the ion thermal velocity.)

Solving the Poisson equation leads to a double Yukawa repulsive potential:

$$\Phi(r) = 1/r \cdot (Z_{SR}^* \exp(-r/\lambda_{SR}) + Z_{LR}^* \exp(-r/\lambda_{LR})) \quad (9)$$

Simulation:

- $\Lambda = \lambda_{LR}/\lambda_{SR}$  as measure for dominance of LR over SR interaction
- demixing is accompanied by domain growth  $L(t) \propto t^\alpha$



# Simulated demixing of a binary complex plasma

[pk3-plus]

[simulation]

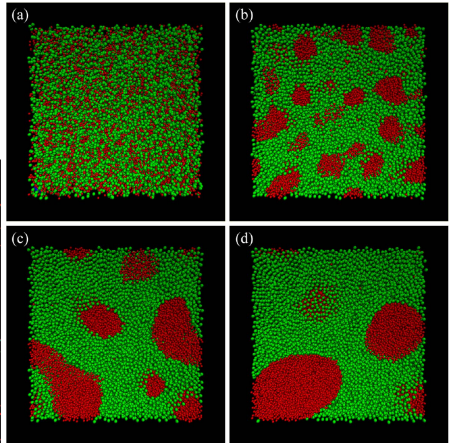
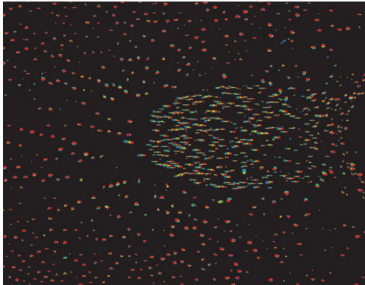


Fig. 5 : Above: Demixing of a binary complex plasma. Ivlev,2009 and Jiang,2011.



# Simulated demixing of a binary complex plasma

Histogram method to get the number of particles in a demixed domain:

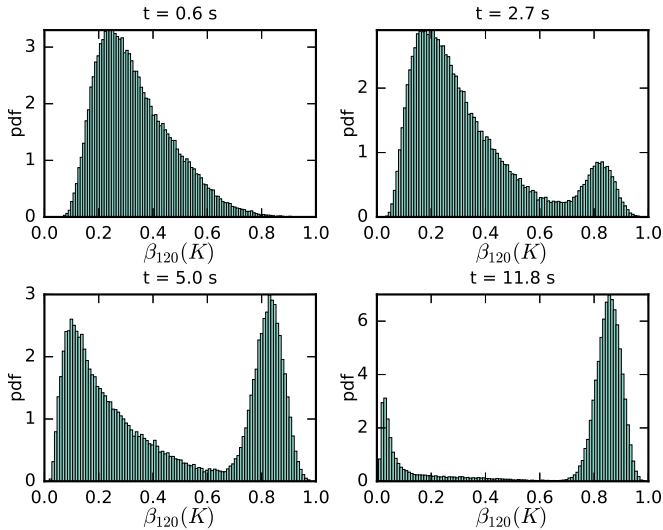


Fig. 6: Above: Growth of minority phase domains for simulated demixing of a complex plasma.



# Simulated demixing of a binary complex plasma

Domain growth for measures  $m \in \{MT0, MT2, \text{power spectrum method}\}$ .

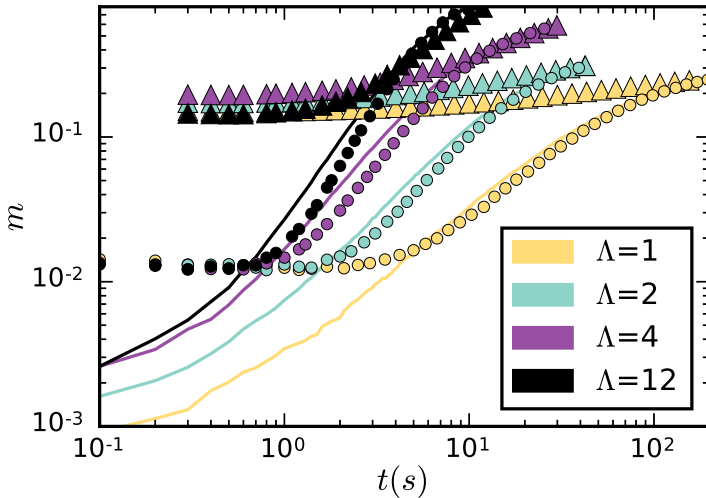


Fig. 7 : Above: Growth of minority phase domains for simulated demixing of a complex plasma. Lines show MT0 analysis, circles MT2 analysis and triangles a power spectrum method.



# Simulated demixing of a binary complex plasma

Slopes of MT measures hints towards universal behavior:

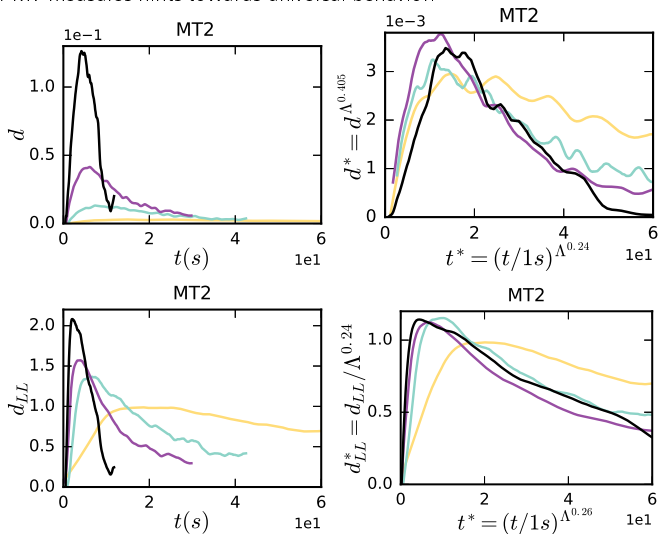


Fig. 8 : Above: Universal behaviour of gradients of MT measures for  $\Lambda \neq 1$ .



# Simulated demixing of a binary complex plasma

Demixing occurs in two stages:

- agglomeration of neighbouring particles
- cascades of merging domains (only for  $\Lambda > 1$ )

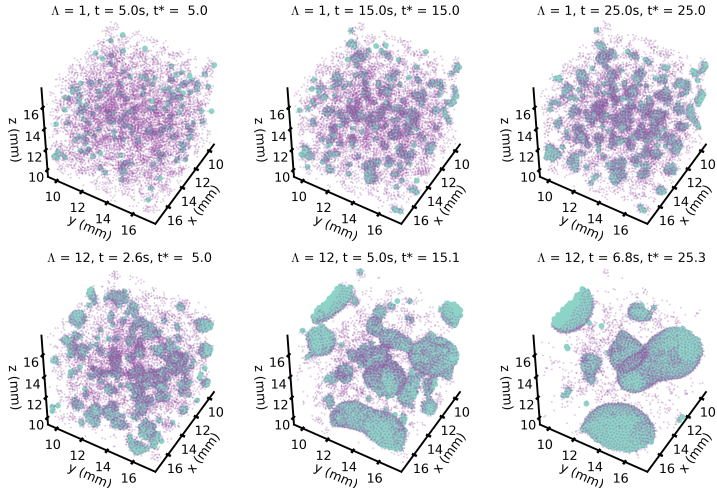


Fig. 9 : Above: Qualitative differences in demixing behavior for  $\Lambda > 1$ .



# Outlook

- investigate transition close to  $\Lambda = 1$
- MT demixing analysis of well know Lennard Jones system
- demixing for general interaction potentials, universality?



**Thank you for your attention!**

